

RESEARCH ARTICLE

In Situ Single-Atom Decoration of Transition Metal Dichalcogenides by Millisecond Flash Thermal Synthesis Toward High Performance Room Temperature Chemiresistors

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ABSTRACT

Transition metal dichalcogenides (TMD) are attractive for adsorption-driven reactions, yet their activity is strongly site-dependent. Activity is concentrated at edge motifs, whereas most exposed area resides on inert basal planes. Dual-site designs that enrich edges while activating basal planes with catalysts remain challenging to realize without coarsening and metal aggregation during thermal processing. Here, an intense pulsed light-driven flash thermal synthesis route is demonstrated that directly converts ammonium tetrathiomolybdate ((NH₄)₂MoS₄) into few-layer, edge-rich MoS₂ nanoflakes (FTS-MoS₂) within 10 ms pulse in ambient-air. Ultrafast photothermal shock (1192–1811°C; ~10⁵/10⁴°C s⁻¹ heating/cooling rates) suppresses in-plane coarsening and out-of-plane stacking, while Pt, Ir, or Au single atoms are uniformly anchored on MoS₂ via rapid metal-sulfur coordination without aggregation. As a proof-of-concept, FTS-MoS₂ exhibits a 23.6-fold higher NO₂ response at 5 ppm than solvothermally synthesized MoS₂. Pt single atom functionalization (FTS-Pt_{SA}-MoS₂, 1.2 wt%) further boosts the response by 22.8-fold versus pristine FTS-MoS₂ and achieves 100.8% response toward 400 ppb NO₂ at room temperature. Density functional theory supports enhanced NO₂ adsorption and charge transfer on FTS-Pt_{SA}-MoS₂. With a low electrical energy input (8.6 kJ g⁻¹) and scalable irradiation, ultrafast FTS enables industrially relevant active-site and single-atom engineering in TMDs for high-performance gas sensors.

1 | Introduction

Transition metal dichalcogenides (TMDs) have emerged as versatile low-dimensional materials whose electronic structure and surface chemistry can be engineered for catalysis, sensing, and

energy conversion [1]. Among them, molybdenum disulfide (MoS₂) is particularly attractive because its surface can host a wide range of adsorption and charge-transfer processes while maintaining chemical robustness [2, 3]. However, a key materials-level principle is that MoS₂ activity is strongly site-dependent.

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Edge sites are generally far more reactive than basal planes, and thermodynamics often favors basal plane exposure during growth, limiting the density of accessible active sites [4, 5]. As a result, downsizing MoS₂ into edge-rich nanoflakes is a direct and powerful strategy to maximize active-site density per mass and to accelerate surface reaction kinetics.

Despite the clear design purpose, controlling MoS₂ at the nanoscale remains challenging. Many conventional synthesis routes (e.g., solvothermal growth, postsulfurization, and furnace-based conversions) require long reaction times and extended thermal budgets, which commonly promote coarsening and structural reorganization [6]. In situ microscopy studies have shown that MoS₂ can grow and enlarge through temperature-driven processes such as merging and ripening-like evolution, ultimately increasing lateral size and reducing the ratio of edge exposure [7]. Consequently, achieving nano-lateral size is still challenging, especially when the synthesis must be performed under prolonged annealing. Although edge enrichment enhances activity, integrating a basal-plane activation strategy further enhances catalytic activity, including in the presence of vacancies and additives. Functionalization via the incorporation of catalytic species, particularly single-atom catalysts (SACs), offers a powerful strategy to boost basal plane activity with maximum atomic efficiency. However, stabilizing and retaining isolated metal atoms on 2D supports is intrinsically difficult because metal species can migrate and aggregate, especially under thermal treatment [8]. Conventional synthesis routes for SACs on MoS₂ often involve prolonged high-temperature treatments under reducing conditions, which can induce lateral domain growth, multilayer formation, or agglomeration of metal atoms, ultimately challenging the synthesis for nanoscale integrity of the MoS₂ host [9]. Therefore, achieving in situ anchoring of SACs while maintaining the nanoscale morphology of MoS₂ is critical for the realization of highly active TMD-based functional materials.

Ultrafast thermal treatment techniques have emerged as promising strategies for nanoscale catalyst synthesis by enabling rapid nucleation while suppressing undesired growth [10]. Among ultrafast approaches, electromagnetic irradiation-based annealing (e.g., microwaves, lasers, and broadband light) is particularly attractive due to noncontact and integration of diverse substrates. However, each modality presents distinct trade-offs. Microwave processing can provide volumetric heating, yet the temperature profile is often sensitive to cavity field distributions [11]. Laser annealing delivers extremely high local intensities but typically relies on beam scanning and overlap strategies to achieve large-area processing [12]. By contrast, xenon-flash/intense pulsed light (IPL) annealing offers millisecond-duration broadband irradiation with high radiant energy density and scalable, large-area exposure, enabling strong photothermal heating under short processing durations [13]. Such flash-type thermal shock processing has enabled the synthesis of single-atom catalysts and high-entropy nanomaterials by confining high-temperature exposure to millisecond durations [14–17]. Although flash thermal synthesis (FTS) process has been exploited for carbon and oxide materials, its extension to TMDs has thus far primarily focused on ultrafast synthesis and phase/heterostructure engineering, whereas catalytic functionalization, particularly the integration of atomically dispersed catalytic sites, remains rarely explored [18, 19].

Here, we demonstrate single-step FTS strategy (1192°C–1811°C peak temperature, 10 ms irradiation, and 10⁴–10⁵°C s^{−1} ramping rates) that directly converts ammonium tetrathiomolybdate ((NH₄)₂MoS₄) into edge-rich MoS₂ nanoflakes, while simultaneously enabling the uniform dispersion of noble-metal single atoms (e.g., Pt, Ir, and Au) on the MoS₂ framework under ambient air. The ultrafast thermal kinetics facilitate the formation of MoS₂ with a high edge-to-volume ratio and promote rapid metal-sulfur coordination of isolated atoms. To demonstrate the practical feasibility of this approach, we investigated chemiresistive gas sensing performance of the FTS-based MoS₂ (FTS-MoS₂) and single atom functionalized MoS₂ (FTS-M_{SA}-MoS₂). Notably, FTS-MoS₂ shows a 23.6-fold higher NO₂ response than solvothermally synthesized MoS₂. Pt single atom functionalization (1.2 wt%) further enhances the response by 22.8-fold at room temperature at 400 ppb NO₂, achieving a limit of detection (LOD) of 17.2 ppb. The highly active sites for gas molecules on TMDs can be broadly expanded to various electrocatalysis fields by diversifying single atoms and extending to other TMD materials, as these properties significantly enhance gas adsorption and desorption characteristics crucial for catalytic performance. In addition, this rapid TMDs synthesis method could realize low-cost, mass production of nanocrystalline TMDs at only 8.6 kJ g^{−1} in electrical energy.

2 | Results and Discussion

2.1 | Flash Thermal Synthesis of TMDs

Figure 1a schematically illustrates IPL-driven FTS of TMDs. First, 100 mg of ammonium tetrathiomolybdate (ATTM, (NH₄)₂MoS₄) was uniformly dispersed in 1 mL of deionized water through ultrasonication. Then, 100 µL of the ATTM solution was drop-cast onto a microscope glass slide and dried on a hot plate at 80 °C (Figure S1). The xenon lamp used in the intense pulsed light (IPL) process emits light over a broad wavelength range of approximately 300–1000 nm (Figure S2). Within this range, ATTM exhibits strong visible light absorption from 400 to 800 nm, as confirmed by the absorption spectra in Figure S3. The absorbed energy was estimated using photon counts (Figure S2) and absorbance (Figure S3). Notably, visible light accounted for 78.7% of the total irradiation during the IPL process, as indicated in Figure S4. This high optical absorbance enables efficient photothermal heating via nonradiative relaxation [20, 21]. After the IPL process, the samples were purified by centrifugation at 8000 rpm for 10 min with deionized water and ethanol. According to previous reports, ATTM thermolysis proceeds via a MoS₃ intermediate with the release of NH₃ and H₂S, followed by conversion to MoS₂ at higher temperatures [22]. At sufficiently high temperatures, sulfur loss/sublimation can generate S vacancies in MoS₂ [23]. After IPL irradiation, visible redistribution of the reacted material was observed on the substrate (Figure S5). Together with the decreasing recovered solid mass before washing at higher FTS powers (Note S1), this observation suggests that additional process-associated loss occurs during the FTS pulse, potentially through volatilization and physical material ejection. Figure 1b shows temperature-time profiles under a single 10 ms IPL pulse at lamp power settings of 1000, 1500, and 2000 W. The resulting peak temperatures reached approximately 1192°C, 1481°C, and 1811°C, respectively, depending on the irradiation

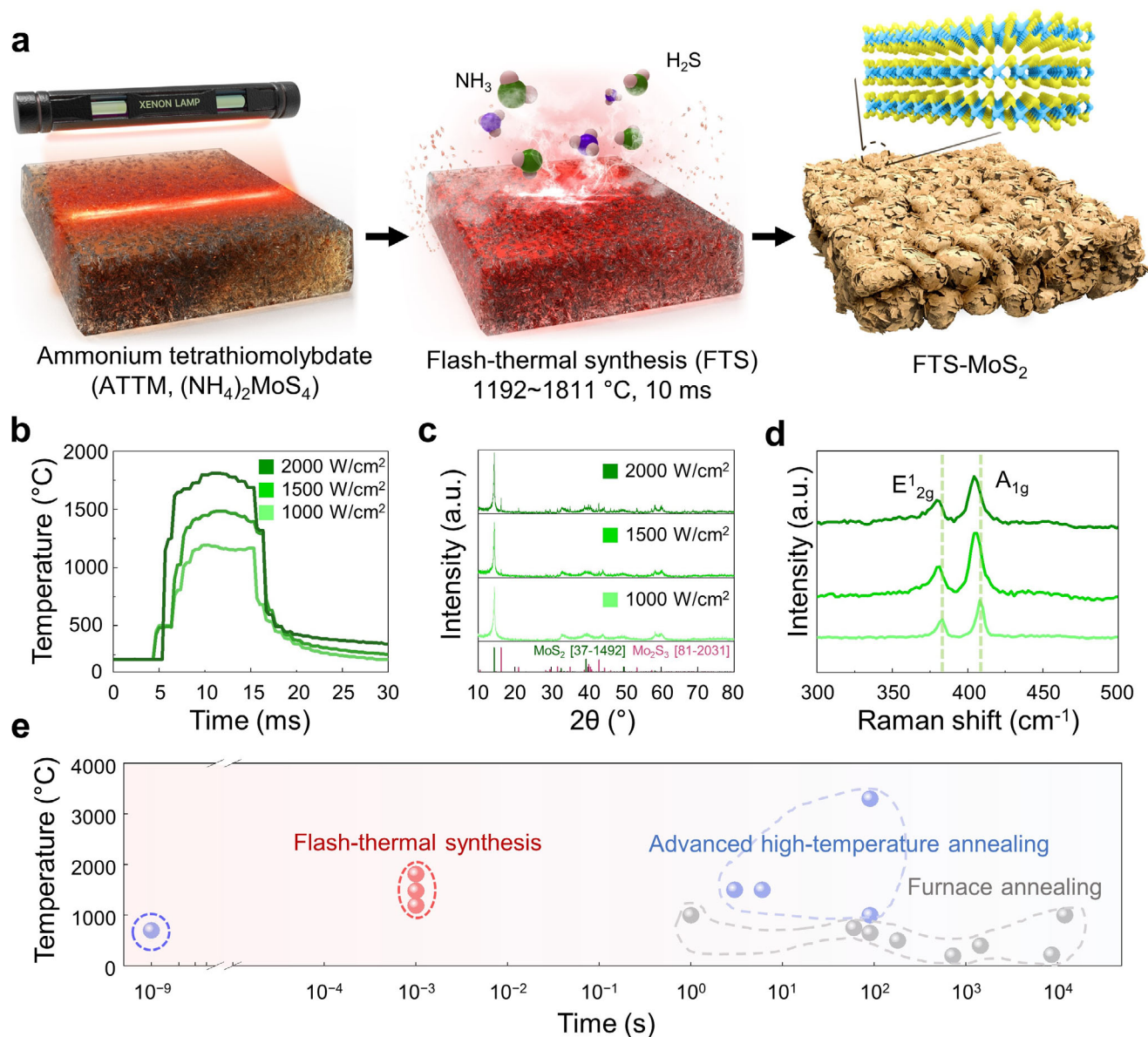


FIGURE 1 | (a) Schematic illustration of the synthesis of MoS₂ via flash thermal synthesis (FTS). (b) Photothermal temperature-time profiles, (c) XRD patterns, and (d) Raman spectra of FTS-MoS₂ synthesized under power densities of 1000, 1500, and 2000 W cm⁻². (e) Comparison of the synthesis time and temperature between FTS, advanced high-temperature annealing, and conventional furnace annealing. The detailed temperature and time parameters are summarized in Table S1.

energy. FTS-based ultrafast heating/cooling of the ATTM coating ($\sim 10^5$ °C s⁻¹ heating rate; $\sim 10^4$ °C s⁻¹ cooling rate; Figure S6) kinetically suppresses excessive growth, enabling few-layer MoS₂ nanoflakes (FTS-MoS₂). Scanning electron microscopy (SEM) images show spherical secondary particles assembled from small MoS₂ nanoflakes (Figures S7 and S8). These structures are attributed to the ultrafast reaction kinetics of the FTS process, which suppresses the lateral growth of MoS₂. In addition, transmission electron microscopy (TEM) was employed to investigate the structural characteristics of FTS-MoS₂, particularly its high edge-to-volume ratio, predominantly short clusters characterized by discontinuous, bent, disjointed, and curved basal planes, indicative of a highly defective structure (Figure S9). In contrast to the solvothermally synthesized MoS₂ (Figure S10), this morphology is consistent with millisecond-scale thermal shock kinetics,

where the extremely short high-temperature annealing time limits diffusion-driven coarsening and interrupts lateral grain growth/coalescence, resulting in fragmented nanocrystalline domains [20]. These short and fragmented flakes suggest a high density of catalytically active edge sites. The high photothermal temperatures achieved under IPL irradiation were sufficient to decompose ATTM, resulting in the formation of crystalline MoS₂, as evidenced by distinct XRD patterns matching the standard PDF#37-1492 (Figure 1c). The crystallinity improved progressively with increasing irradiation energy, which can be attributed to the corresponding increase in FTS temperature (Figure S11). In addition, as the photothermal temperature increased, sulfur-deficient conditions generated during ATTM decomposition promoted the formation of a Mo₂S₃ phase (PDF#81-2031), which the ultrafast cooling process could kinetically retain [24]. Notably, although

the XRD patterns indicate the presence of this characteristic Mo_2S_3 phase, the XPS spectra are dominated by Mo^{4+} ($\text{Mo } 3d_{5/2}$ at $\sim 229.2\text{--}229.5$ eV) and S^{2-} ($\text{S } 2p$ at $\sim 162.0\text{--}162.3$ eV), together with a minor MoO_x contribution arising from surface oxidation (Figure S12) [25]. The absence of crystalline MoO_x peaks in XRD suggests that this oxidized component is surface-limited and/or amorphous rather than a bulk oxide phase. This surface-limited oxidation is due to the millisecond-scale thermal profile, in which the 10 ms of processing followed by rapid quenching restricts oxygen diffusion and suppresses extensive bulk oxidation (Figure S13) [26]. The surface of the resulting materials is predominantly composed of MoS_2 species with mild oxidation, whereas the Mo_2S_3 phase is more located in the inner sites in the materials and therefore is detected by XRD rather than by XPS. In this regard, the ultrafast FTS process generates a temperature gradient across the precursor domain, such that the hotter interior region favors the formation of the sulfur-deficient Mo_2S_3 phase, while the outer region is converted to MoS_2 and partially oxidized upon exposure to ambient air. Such spatially heterogeneous phase evolution is also consistent with previous flash-heating studies on ATTM/carbon systems, where an interior of high-temperature-derived $\alpha\text{-MoC}$ and an exterior of relatively low-temperature-derived MoS_2 were formed under FTS, giving rise to a core-shell heterostructure [27]. Raman spectroscopy further confirmed the formation of 2H- MoS_2 (Figure 1d). Two characteristic peaks at 383 and 408 cm^{-1} were observed, corresponding to the in-plane (E_{1g}) and out-of-plane (A_{1g}) vibrational modes of 2H- MoS_2 , respectively. As the irradiation energy increased, the E_{1g} peak shifted to 381 cm^{-1} and the A_{1g} peak to 404 cm^{-1} , indicating a redshift in both vibrational modes due to lattice thermal expansion and anharmonic effects owing to sulfur vacancy formations. In addition, the weak but discernible peaks at 227 and 454 cm^{-1} , corresponding to the LA(M) and 2LA(M) modes, respectively, further indicate the defect-rich nature of FTS- MoS_2 (Figure S14). These sulfur vacancies became more pronounced at higher photothermal temperatures under intense irradiation [28]. X-ray photoelectron spectroscopy (XPS) analysis further suggested progressive sulfur deficiency with increasing irradiation energy, as reflected by a decrease in the S:Mo ratio from 1.996 (1000 W) to 1.990 (1500 W) and 1.920 (2000 W) (Figure S15) [29]. In addition, atomic resolution STEM line profiles reveal a systematic contraction of the nearest-neighbor Mo–Mo spacing (mean of 40–50 measured distances), decreasing from 3.22 to 3.12 and 2.96 Å as the irradiation energy increased (Figure S16). This trend is consistent with local lattice relaxation/strain accumulation around sulfur vacancies, where adjacent Mo atoms relax inward upon S removal [30, 31]. DFT calculations show that introducing sulfur vacancies leads to a reduced Mo–Mo distance, reproducing the experimentally observed contraction and supporting vacancy-induced lattice strain as the origin of the measured Mo–Mo spacing change (Figure S17).

Altogether, this FTS route offers access to high temperatures and faster ramping rates than many reported TMD synthesis/annealing approaches (Figure 1e and Table S1). Conventional furnace heating typically operates at temperatures around 1000°C and requires several hours to synthesize TMDs where large amounts of energy are used to maintain the temperature of the chamber. Although laser annealing has been reported to reach temperatures of $\sim 700^\circ\text{C}$ within 2 ns, the processed area (spot size) is limited to 20 μm [32]. In comparison, the FTS method

rapidly achieves temperatures up to 1811°C within just 10 ms while providing a significantly larger irradiation area of up to 75 cm^2 , determined by the size of the xenon lamp (5×15 cm in this system). Owing to ultrafast ramping, direct precursor-to-product conversion, and millisecond/large processing, the FTS-based synthetic method is energy-efficient (8.6–15.5 kJ g^{-1}) and enables MoS_2 synthesis in ambient air (Notes S1 and S2).

2.2 | Flash Thermal Synthesis of SAC-TMDs

Among the irradiation conditions, we selected the 2000 W setting as an optimized FTS condition because it provides a sufficiently high temperature ($> 1700^\circ\text{C}$) to drive the synthesis of defective MoS_2 nanoflakes and simultaneously decomposition of noble metal precursors within a single 10 ms pulse (Figure 2a). Importantly, the millisecond-duration annealing and rapid quenching can kinetically suppress metal aggregation, enabling atomically dispersed metal species to form on the basal planes of MoS_2 . Given that isolated metal atoms can be stabilized in different coordination environments on the MoS_2 surface, we selected representative noble metals to demonstrate the versatility of the FTS-based SAC functionalization strategy [33, 34]. Noble metal precursors, such as $\text{H}_2\text{PtCl}_6 \cdot 6\text{H}_2\text{O}$, IrCl_3 , and $\text{HAuCl}_4 \cdot 3\text{H}_2\text{O}$ were introduced into the ATTM solution to achieve one-step SAC functionalization during FTS. To remove residual precursor species, the samples were thoroughly washed with ethanol prior to characterization (Figure S18). To investigate the energetically favorable adsorption site of the Pt, Ir, and Au single atom on the basal plane, density functional theory (DFT) calculations were conducted. To model the basal plane of FTS- MoS_2 , the slab model with (001) plane was constructed from the bulk 2H- MoS_2 structure, including Mo_2S_3 and sulfur vacancies (Figure S19). Unique adsorption sites were defined on the surface of the FTS- MoS_2 slab model (Figure 2b), consisting of Mo_{near} , Mo_{mid} , and Mo_{far} site where the single atom is located on the top of Mo atoms near, at midrange, and far from the sulfur vacancy, respectively, and a S site where the single atom is located on a top of S atom. Figure 2c shows the adsorption energies of Pt, Ir, and Au on the defined adsorption sites. The energetically most stable structures for every single atom are shown in Figure 2d. For Pt and Ir single atom, the Mo_{mid} site is preferred adsorption site with a large difference of adsorption energy (0.68 and 1.51 eV for Pt and Ir, respectively). However, the Au single atom prefers the S site with a relatively small energy difference (0.24 eV), indicating a relatively high probability of occupying other adsorption sites as well (Figure S20). Figure 2e–g provide microscopic and spectroscopic evidence for atomically dispersed Pt on MoS_2 (FTS-Pt_{SAC}- MoS_2 , Figure S21). In the Pt 4f region, the Pt 4f_{7/2} component at 73.09 eV is shifted to higher binding energy relative to metallic Pt (71 eV), consistent with cationic Pt species ($\text{Pt}^{\delta+}$) interacting with sulfur in the MoS_2 (Figure S22) [35]. Bright contrast spots in the high-angle annular dark-field scanning transmission electron microscopy (HAADF-STEM) images were attributed to Pt single atoms, which were further verified by line scan profiling (Figure 2f) and 3D contour mapping (Figure 2g). Similarly, the formation of Ir (Figure 2h–j and Figure S23) and Au (Figure 2k–m and Figure S24) single atoms on MoS_2 was confirmed. In the XPS analysis, the peaks of Ir and Au were located at 62.22 and 86.02 eV for FTS-Ir_{SAC}- MoS_2 and FTS-Au_{SAC}- MoS_2 , respectively, corresponding to $\text{Ir}^{\delta+}$ and $\text{Au}^{\delta+}$ states. These peaks

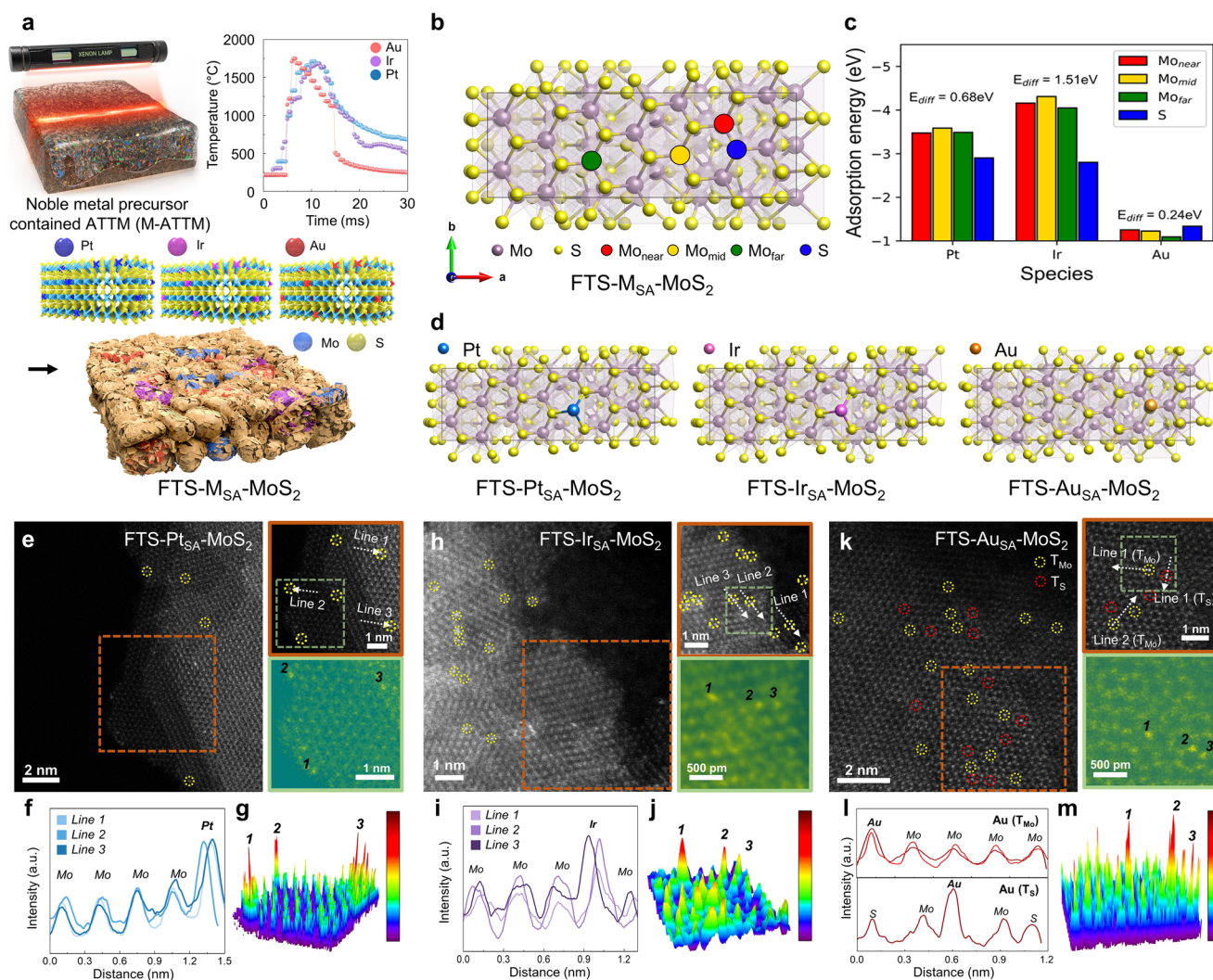


FIGURE 2 | (a) Schematic illustration of metal-containing ATTMs (M-ATTMs) to produce single atom functionalized MoS₂ via FTS (FTS-M_{SA}-MoS₂; M = Pt, Ir, and Au) and photothermal temperature-time curves. (b) FTS-MoS₂ slab model and adsorption sites considered in DFT calculations (c) Calculated adsorption energies of Pt, Ir, and Au at top of Mo_{near}, Mo_{mid}, Mo_{far}, and S sites. (d) DFT-optimized configurations of Pt, Ir, and Au single atoms on MoS₂ at their preferred adsorption sites. HAADF-STEM images and their corresponding line profile, and 3D-contour mapping images for (e-g) FTS-Pt_{SA}-MoS₂, (h-j) FTS-Ir_{SA}-MoS₂, and (k-m) FTS-Au_{SA}-MoS₂.

suggest the formation of strong metal-chalcogen interactions (Figures S25 and S26) [36, 37]. The formation of isolated atoms without detectable metallic properties is attributed to the ultrafast, nonequilibrium characteristics of the FTS process, which enables rapid precursor conversion/anchoring while kinetically suppressing coarsening that typically drives nanoparticle (NP) formation. High-power FTS (2000 W) led to sulfur deficiency at the MoS₂ surface (4%–8% by XPS, Figures S15 and S27). Notably, high-temperature desulfurization of MoS₂ is known to induce sulfur-deficient reconstructions [38]. Moreover, prior studies indicate that introducing isolated metal atoms can promote defect generation and structural/phase evolution in MoS₂, which may facilitate sulfur loss under ultrafast thermal shocks [39]. The excessive sulfur vacancies can lead to a phase change from MoS₂ to Mo₂S₃, especially in metastable conditions in ultrafast synthesis, as confirmed by XRD analysis (Figure S28), similar to the previous reports [27]. The Pt-, Ir-, and Au-functionalized FTS-MoS₂ exhibited no detectable metallic peaks, confirming that the noble metal species were atomically dispersed without the

formation of metallic NPs (FTS-M_{SA}-MoS₂, M denotes Pt, Ir, and Au).

To identify the Pt loading regime that yields atomically dispersed species on FTS-MoS₂, we synthesized a series of FTS-Pt-MoS₂ with different nominal Pt contents. ICP-MS confirmed Pt loadings of 0.03, 0.5, 0.7, 1.2, and 3 wt%. The XRD analysis indicated the MoS₂ and Mo₂S₃ with all FTS-Pt-MoS₂ and exhibited no detectable metallic Pt peaks, which is consistent with the absence of Pt NPs and/or the detection limit of ultrasmall/low-loading Pt by XRD. (Figure 3a). Notably, the relative Mo₂S₃ contribution increased with increasing Pt loading, suggesting that Pt could facilitate sulfur depletion during the high-temperature FTS process and thereby promote the formation of sulfur-deficient Mo₂S₃. In addition, the electron spin resonance spectrum showed a strong signal at a g value of 2.21, indicating the presence of defect-related paramagnetic centers that may be associated with a sulfur-deficient local Mo-S environment, including sulfur-vacancy sites (Figure S29) [40]. Raman spectra further reveal

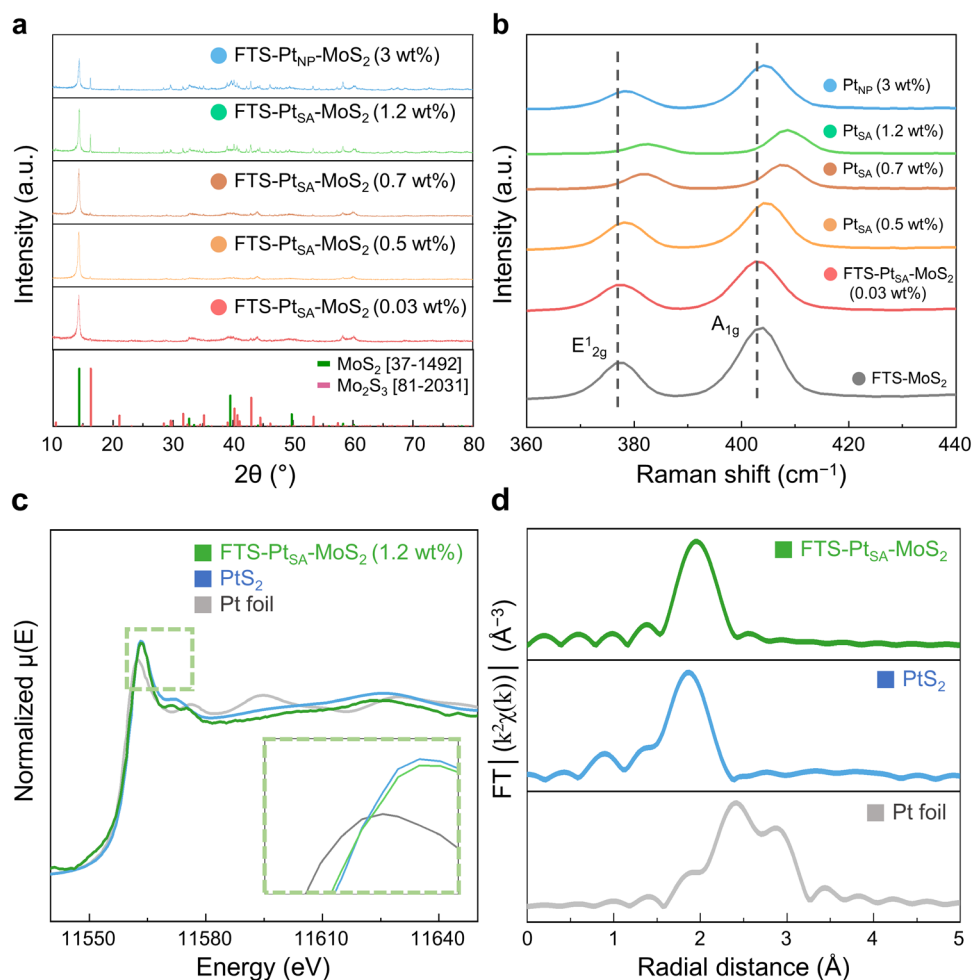


FIGURE 3 | (a) XRD patterns of FTS-Pt-MoS₂ synthesized with different Pt loadings (0.03, 0.5, 0.7, 1.2, and 3 wt%). (b) Raman spectra of FTS-MoS₂ and FTS-Pt-MoS₂. (c) Pt L₃-edge XANES spectra of FTS-Pt_{SA}-MoS₂ (1.2 wt%) compared with PtS₂ and Pt foil. (d) Fourier-transformed k²-weighted EXAFS magnitude of FTS-Pt_{SA}-MoS₂ (1.2 wt%) with PtS₂ and Pt foil.

a Pt loading dependent shift of the E_{12g}¹ and A_{1g} modes. The low loading (0.03 wt%) induces negligible changes, whereas intermediate loadings in the single atoms show blue shifts, consistent with charge-transfer/strain effects that are known to modulate MoS₂ phonon frequencies (Figure 3b) [41, 42]. At the highest Pt loading investigated (3 wt%), the Raman-shift trend was reversed, coinciding with the emergence of a metallic Pt component in XPS and Pt-containing nanoparticles in electron microscopy (Figures S30 and S31). These observations indicate a loading-dependent loss of atomic dispersion under the present FTS conditions. The formation of aggregated Pt species at 3 wt% may reflect an increased probability of Pt–Pt association relative to isolated stabilization as the surface Pt population increases. However, because the density and distribution of the available anchoring sites on FTS-MoS₂ were not independently quantified, a specific anchoring-capacity threshold cannot be assigned from the present data. Accordingly, the 1.2 wt% sample, which retained predominantly atomically dispersed Pt under the present FTS conditions, was selected as the representative sample for detailed structural characterization. X-ray absorption near-edge spectroscopy (XANES) and extended x-ray absorption fine structure (EXAFS) profiles were used to investigate the

electronic structure and coordination of the Pt in FTS-Pt_{SA}-MoS₂ (1.2 wt%). The XANES analysis on FTS-Pt_{SA}-MoS₂ (1.2 wt%) shows that their Pt L₃ edge absorption energy values and White line intensities are significantly higher than that of metallic Pt foil, which indicating the Pt species in the FTS-Pt_{SA}-MoS₂ are dominantly in the oxidized state (Figures 3c and S32). These XANES results are consistent with the XPS analysis, further confirming that the Pt species in FTS-Pt_{SA}-MoS₂ are predominantly oxidized rather than metallic (Figure S33). The Fourier-transformed EXAFS of FTS-Pt_{SA}-MoS₂ exhibits a dominant peak at 1.9 Å, which can be attributed to the Pt–S scattering (Figure 3d). In addition, the corresponding EXAFS fitting results in both k-space and R-space, together with the Wavelet-transformed EXAFS contour analysis, further support a sulfur-coordinated local environment around Pt (Figure S34). The fitting results demonstrate that the Pt–S bond has an average bond length of 2.36 Å and an average coordination number of 3.2 in the FTS-Pt_{SA}-MoS₂ (Table S2). The first-shell Pt–S feature of FTS-Pt_{SA}-MoS₂ appears at a slightly higher R value than that of PtS₂, which may be attributed to the undercoordinated and locally distorted Pt–S environment of atomically dispersed Pt species [43].

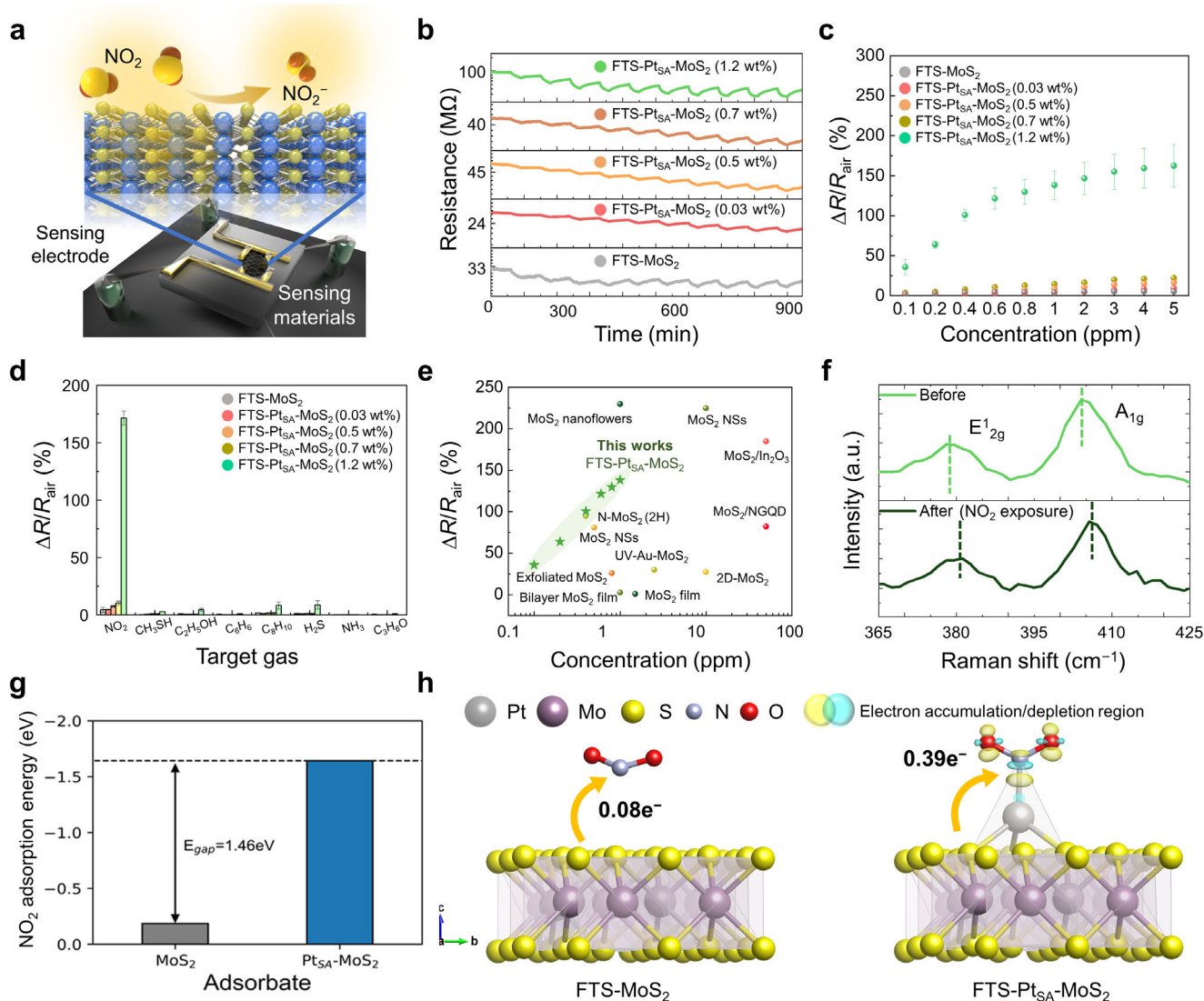


FIGURE 4 | (a) Schematic illustration of the Al₂O₃-based Au interdigitated-electrode device and the proposed charge-transfer sensing mechanism of NO₂ on FTS-MoS₂. (b) Dynamic resistance traces for FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ under NO₂ exposure (0.1–5 ppm) at room temperature. (c) Dynamic gas sensing performance of FTS-MoS₂ and FTS-Pt_{SA}-MoS₂. (d) Selective sensing performances of FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ toward 5 ppm gases. (e) Comparison of the NO₂ response values to the state-of-the-art MoS₂-based sensing materials. (f) Ex situ Raman spectra for FTS-Pt_{SA}-MoS₂ (1.2 wt%) before and after NO₂ exposure. (g) DFT-calculated adsorption energies of NO₂ on bare MoS₂ and Pt_{SA}-MoS₂. (h) Charge density difference for NO₂ adsorption on bare MoS₂ and Pt_{SA}-MoS₂, illustrating enhanced charge transfer mediated by Pt single atoms.

2.3 | Applications: Chemiresistive NO₂ Gas Sensors

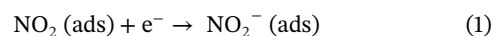
To demonstrate the practical utility of FTS-derived MoS₂, chemiresistive NO₂ sensing performance was evaluated. From an environmental and occupational-health perspective, sensitive detection of NO₂ at low-ppm levels is important; notably, the United States Occupational Safety and Health Administration (OSHA) lists a ceiling permissible exposure limit of 5 ppm for NO₂. Accordingly, the development of highly sensitive sensing materials is crucial for accurate detection and regulatory compliance in environmental and occupational settings. Homemade gas sensors were fabricated by drop-coating of sensing material solutions onto custom Al₂O₃ substrates. The alumina substrates were patterned with interdigitated Au electrodes, featuring a width of 25 μm and a distance of 70 μm (Figure 4a). Since

gas adsorption on MoS₂ is strongly site-dependent, edge-/defect-rich MoS₂ architectures typically exhibit higher reactivity than basal plane-dominant counterparts [4]. In addition, the sulfur vacancies serve as active sites for gas adsorption, further improving sensing performance [44]. Consistent with this concept, FTS-MoS₂, which has smaller lateral features and higher edge density than solvothermally synthesized MoS₂, exhibited a 23.6-fold higher response to 5 ppm NO₂, indicating that edge-rich nanoflakes effectively improve gas-surface interactions (Figures S9 and S35). In addition, the heterointerfaces, such as Mo₂S₃ and MoO_x, in the FTS-MoS₂ also contribute to the enhanced response toward oxidizing NO₂ through interfacial charge transfer and hole accumulation [45]. Furthermore, catalyst functionalization significantly contributed to the enhancement of sensing performance, underscoring its pivotal role in boosting both sensitivity and selectivity [46]. The gas sensing characteristics of

Pt-functionalized FTS-MoS₂ (FTS-Pt-MoS₂) were systematically compared to those of FTS-MoS₂, further highlighting the catalytic role of Pt in enhancing sensor performance (Figure 4b). FTS-MoS₂ exhibited a p-type-like sensing response, as indicated by a decrease in resistance upon NO₂ exposure and positive apparent Hall slopes obtained from fixed-contact Hall-effect measurements (Figures S36, S37, and Table S3). We note that this macroscopic response should not be assigned to a single microscopic origin. FTS-MoS₂ contains sulfur deficiency, surface-limited MoO_x species, MoS₂/Mo₂S₃ phase heterogeneity, and multiple interfacial environments, all of which can perturb carrier density and charge transport and may collectively contribute to the observed p-type-like behavior. FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ sensors, with varying Pt loadings (0.03, 0.5, 0.7, and 1.2 wt%), were exposed to NO₂ at different concentrations from 0.1 to 5 ppm under dry conditions (Figure 4c). The responses 5.38%, 9.28%, 16.35%, 22.07%, and 162.44% were achieved from FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ (0.03, 0.5, 0.7, and 1.2 wt%) to 5 ppm NO₂, respectively. Among them, FTS-Pt_{SA}-MoS₂ (1.2 wt%) exhibited the highest sensing performance due to its maximized catalytic effect of single Pt atoms. Notably, outstanding sensing performance was maintained down to low ppb-level concentrations, with a 100.8% response measured at 400 ppb NO₂, 22.8-fold higher response than FTS-MoS₂. At a higher Pt loading of 3 wt%, the excessive catalyst could hinder interactions with target gas molecules, resulting in a reduced sensing response (Figure S38). The FTS-MoS₂ exhibited excellent selectivity toward 5 ppm NO₂ when compared to other analytes at the same concentration, including CH₃COCH₃, NH₃, C₇H₈, C₈H₁₀, C₆H₆, and C₂H₅OH (Figures 4d and S39). To support the selectivity results shown in Figure 4d, additional DFT calculations of the adsorption energies and Bader charge differences were conducted for seven gas molecules (CH₃SH, C₂H₅OH, C₆H₆, C₈H₁₀, H₂S, NH₃, and C₃H₆O) to compare these results with those of NO₂. Two adsorption configurations were considered for each molecule (Figure S40). As shown in Figure S41, for the most stable adsorption structures, the adsorption energies of each molecule were calculated for CH₃SH (−1.69 eV), C₂H₅OH (−0.92 eV), C₆H₆ (−1.46 eV), C₈H₁₀ (−1.54 eV), H₂S (−1.40 eV), NH₃ (−1.32 eV), and C₃H₆O (−0.99 eV). These values are comparable to those of NO₂ (−1.64 eV) or, in some cases, stronger by up to 0.63 eV. Figure S42 shows the Bader charge differences of the gas molecules on FTS-Pt_{SA}-MoS₂. Among the eight adsorbates, CH₃SH (−0.25 e[−]), C₂H₅OH (−0.11 e[−]), C₆H₆ (−0.04 e[−]), C₈H₁₀ (−0.08 e[−]), H₂S (−0.18 e[−]), NH₃ (−0.18 e[−]), and C₃H₆O (−0.07) exhibit negative charge differences, indicating that these gas molecules are oxidized upon adsorption. In contrast, NO₂ shows a positive charge difference of 0.39 e[−], indicating electron transfer from FTS-Pt_{SA}-MoS₂ to the adsorbed NO₂ molecule. These calculations indicate that NO₂ exhibits a distinct electron-withdrawing interaction among the examined adsorbates on the FTS-Pt_{SA}-MoS₂ model. This preferential charge-transfer behavior is consistent with the enhanced NO₂ response and selectivity, although it does not by itself establish the microscopic origin of the macroscopic p-type-like transport response. Furthermore, the sensing performance of FTS-MoS₂ was evaluated under varying relative humidity (RH) levels and detection limits, along with an evaluation of long-term operational stability. Although excessive H₂O adsorption can hinder NO₂ sensing by competing for active sites, adsorbed water can also modify the surface charge state of MoS₂ and increase the baseline resistance by acting as an electron-trapping

species [47]. The baseline resistance increased under 50% RH and 75% RH for all samples, while the recovery characteristics of FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ (0.03, 0.5, 0.7, and 1.2 wt%) were clearly improved (Figure S43). Water-assisted surface charge transfer and the partial suppression of direct NO₂ chemisorption can improve the recovery behavior at room temperature. In particular, the response of FTS-Pt_{SA}-MoS₂ (1.2 wt%) increased from 158.23 in dry air to 193.37 at 50% RH and 182.68 at 75% RH, indicating a humidity-assisted enhancement of room-temperature NO₂ sensing under moderate RH conditions (Figure S44). Although the dual-site activation strategy effectively enhances NO₂ adsorption and sensing response, the recovery behavior remains limited at room temperature owing to the intrinsic adsorption–desorption trade-off. These results suggest that additional activation strategies, such as light irradiation, mild thermal input, water-assisted modulation, or further tuning of the intrinsic surface activity, are required to improve reversibility. Notably, the FTS-Pt_{SA}-MoS₂ (1.2 wt%) sensors exhibited a low limit of detection (LOD) of 17.2 ppb (Figure S45). Furthermore, it consistently maintained stable sensitivity even after repeated exposure and recovery cycles (25 cycles) toward 5 ppm NO₂, as shown in Figure S46. In addition, FTS-Pt_{SA}-MoS₂ (1.2 wt%) maintained a robust sensing response under ambient air, exhibiting only 2.6% degradation after 3 weeks (Figure S47). To elucidate the stable sensing performance, post-sensing XRD, XPS, and HAADF-STEM analyses were conducted, confirming the preservation of the MoS₂ phase and atomically dispersed Pt species after 25 sensing cycles, despite a mild increase in surface oxidation (Figures S48–S51). This performance indicated superior sensitivity and durability, making it suitable for practical applications requiring long-term operation and high detection accuracy. Overall, FTS-Pt-MoS₂ demonstrated selective NO₂ detection with improved sensing performance, showing superior performance compared with other MoS₂-based and SAC-based room temperature operating NO₂ sensors reported in the literature (Figure 4e and Tables S4, S5).

To elucidate the sensing mechanism of NO₂ on the FTS-Pt-MoS₂ surface, ex situ Raman spectroscopy was conducted after NO₂ exposure. FTS-MoS₂ exhibited a decrease in resistance upon NO₂ exposure, consistent with the p-type-like chemiresistive response observed under the present measurement conditions (Figure 4b). Because the FTS-derived material contains sulfur deficiency, surface oxidation, MoS₂/Mo₂S₃ phase heterogeneity, and multiple interfaces, the macroscopic resistance response likely reflects the combined contribution of these structural and electronic features rather than a single carrier-generation mechanism. NO₂ acts as an electron-accepting adsorbate, and its interaction with the sensing surface can be represented in a simplified form as:



For FTS-Pt-MoS₂, strongly adsorbed NO₂ acts as an oxidant, extracting electrons from MoS₂ and enhancing charge transfer.

After NO₂ exposure, the Raman modes shifted to higher wavenumbers relative to those before exposure (Figure 4f and Figure S52). This ex situ blueshift indicates an overall exposure-induced perturbation of the MoS₂ vibrational/electronic state and is consistent with electron depletion; however, contributions from strain, defects, and structural heterogeneity cannot be excluded. Importantly, ex situ Raman spectroscopy does not

provide site-specific information and therefore cannot directly identify Pt single atoms as the NO₂ adsorption sites. The specific role of Pt single atoms was instead evaluated using DFT calculations for NO₂ adsorption on bare FTS-MoS₂ and FTS-Pt_{SA}-MoS₂ slab models. To consider the various NO₂ adsorption configurations in each surface model, adsorption energies were evaluated for 3 and 2 adsorption configurations on bare FTS-MoS₂ (Figures S53 and S54) and FTS-Pt_{SA}-MoS₂ (Figures S55 and S56) slab models, respectively. Figure 4g shows the representative adsorption energies of NO₂ for each model, demonstrating the stronger adsorption energy of NO₂ on FTS-Pt_{SA}-MoS₂, with a significant energy difference (1.46 eV) relative to the bare FTS-MoS₂. This calculation supports a preferential NO₂ interaction in the presence of the Pt single atom. To investigate the mechanism of enhanced NO₂ adsorption on Pt single atoms, Bader charge analysis and charge density difference maps were employed (Figure 4h) [48]. In the bare FTS-MoS₂ model, NO₂ adsorption resulted in a charge transfer of 0.08 e⁻ from MoS₂ to NO₂. In contrast, the FTS-Pt_{SA}-MoS₂ model showed a charge transfer of 0.39 e⁻ from the Pt_{SA}-MoS₂ to NO₂. Additionally, charge density difference maps demonstrate that charge transfer between the Pt single atom and NO₂ is significantly larger than that between the MoS₂ and NO₂. Together, these DFT results provide computational support for enhanced NO₂ adsorption and charge transfer associated with Pt single-atom functionalization. The FTS-Ir_{SA}-MoS₂ and FTS-Au_{SA}-MoS₂ showed similar enhancements in NO₂ adsorption and charge redistribution, supporting the enhanced NO₂ selectivity and sensing responses (Figures S57–S59). These results suggest that the FTS strategy is broadly applicable to diverse single-atom functionalization on MoS₂ and that catalytic tuning of the single-atom species can be used to tailor gas adsorption/desorption and sensing behavior.

3 | Conclusion

In summary, we demonstrate that IPL-driven FTS based on millisecond-scale photothermal shock provides a simple, energy-efficient, and scalable route to engineer highly active sites in MoS₂ under ambient air conditions. This process simultaneously suppresses lateral grain growth and yields edge-abundant and sulfur vacancy-rich few-layer MoS₂ nanoflakes. Moreover, when noble-metal precursors are introduced, the same one-step FTS enables in situ anchoring of atomically dispersed Pt, Ir, and Au single atoms without detectable nanoparticle formation up to an optimal loading (e.g., 1.2 wt% Pt). As a proof-of-concept application, FTS-MoS₂ exhibits a 23.6-fold higher NO₂ response than solvothermally synthesized MoS₂, while Pt single atom functionalization (FTS-Pt_{SA}-MoS₂, 1.2 wt%) further boosts the response 22.8-fold versus pristine FTS-MoS₂, achieving a 100.8% response at 400 ppb NO₂, together with a 17.2 ppb detection limit, high selectivity, and robust long-term and humidity-tolerant stability. Ex situ Raman spectroscopy supports an exposure-induced perturbation of the MoS₂ surface, whereas DFT calculations specifically predict stronger NO₂ adsorption and enhanced charge transfer in the presence of Pt single atoms. Collectively, these results establish ultrafast FTS as a powerful platform for active-site engineering in TMDs, offering a general and industrially viable strategy to construct single atom-decorated, defect-rich 2D catalysts for high-performance chemiresistive gas sensors and, more broadly, for catalytic and energy-conversion applications.

4 | Methods

4.1 | Materials

Ammonium tetrathiomolybdate ((NH₄)₂MoS₄, ATTM), HAuCl₄•3H₂O, IrCl₃, H₂PtCl₆•6H₂O, were purchased from Sigma-Aldrich. Chemicals were used without purification.

4.2 | Sample Preparation

To synthesize MoS₂ using the FTS process, we prepared an ATTM solution dispersed in deionized water. Specifically, 100 mg of ATTM was added to 1 mL of deionized water, followed by stirring for 1 h and sonicated for 30 min to disperse uniformly. A 100 μL of ATTM solution was then dropped onto a glass substrate (microslide glass, 76 × 52 × 1.2 mm) using a micropipette and subsequently dried on a hot plate at 80 °C for 2 h. The preparation procedure for noble metal-loaded MoS₂ followed the same steps as for MoS₂, with the exception of adding noble metal precursors (appropriate amount of HAuCl₄•3H₂O, IrCl₃, and H₂PtCl₆•6H₂O) to the ATTM solutions. The coating and drying processes were consistent with the ATTM solution. To synthesize MoS₂ via solvothermal method, 70 mg of ATTM was added to 40 mL of *N,N*-dimethylformamide (DMF) and stirred to achieve a uniform mixture. The resulting solution was transferred to a 50 mL Teflon-lined stainless-steel autoclave and heated at 200 °C for 12 h with a ramping rate of 10 °C/min. After the solvothermal reaction, the black precipitates were collected by centrifugation at 4000 rpm and washed several times with DMF and ethanol to remove residual precursors.

4.3 | IPL System

The intense pulsed light (IPL) system comprises six key components: a simmer board, xenon lamp, capacitors, a DC power supply, an insulated-gate bipolar transistor (IGBT), and an infrared (IR) sensor system (Figure S60). The xenon lamp emits intense pulsed light with a broad wavelength spectrum ranging from 300 to 1000 nm (Figure S2). To initiate light emission, the simmer board ignites the xenon lamp, triggering a streamer breakdown. Following this, the capacitors (charged by the DC power supply to a target voltage (0–500 V)) are discharged through the IGBT, generating an arc breakdown that produces a high-intensity light pulse. The light power density can be precisely controlled by adjusting parameters such as the applied voltage, pulse duration, number of pulses, and the distance between the xenon lamp and the sample. The energy of the emitted pulses was quantified using a laser power meter (NOVA II, Ophir). During the process, the glass substrates were positioned 2–4 cm from the xenon lamp and irradiated for 10 ms at high energy densities of 10, 15, and 20 J cm⁻² (1000, 1500, and 2000 W cm⁻²).

4.4 | Photothermal Temperature of ATTM

The IR sensor systems are introduced to detect the target temperature during the IPL process. The elaborate low temperature detected IR sensor (CTlaser 3MH2, Optris) detects from 200 °C to 1500 °C. The sensor produces a focused dot to measure the sample

temperature, with samples positioned at the proper distance. The measured analog signals are converted to digital signals by a data acquisition board (NI USB-6341 X series Multifunction DAQ, National Instruments). Since the temperature data were obtained with the emissivity value of the IR sensor at 1.00, these data were revised using tabulated values for the emissivity. The following equation (Equation 2) was used for all samples to calculate the revised temperature data.

$$\frac{1}{T_{\text{measured}}} = \frac{1}{T_{\text{corrected}}} + \frac{\lambda \ln \left(\frac{\epsilon_{\text{measured}}}{\epsilon_{\text{corrected}}} \right)}{C_2} \quad (2)$$

In this equation, T_{measured} represents the raw temperature data obtained from the sample, where the emissivity $\epsilon_{\text{measured}}$ is set to 1.00 by the infrared (IR) sensor. $T_{\text{corrected}}$ is the temperature corrected for the actual emissivity. The term λ denotes the detection wavelength of the IR sensor, which is 2.3 μm for the CTlaser. Lastly, C_2 is the second radiant constant, with a value of 14387.8 $\mu\text{m K}^{-1}$. Since the temperature data were obtained with the emissivity set to 1.00, the data were revised using tabulated emissivity values for ammonium tetrathiomolybdate (ATTM). Based on IR camera detection at 2.3 μm , we examined absorbance at 2.3 μm in the FT-IR analysis of the ATTM. For the optically opaque ATTM coating ($T \approx 0$), the spectral emissivity at 2.3 μm was approximated using Kirchhoff's law as $\epsilon \approx \alpha = 1 - R$, where R is the spectral reflectance at 2.3 μm .

4.5 | Material Characterization

Transmission electron microscopy (TEM), scanning transmission electron microscopy (STEM), and elemental analysis (energy-dispersive x-ray spectroscopy (EDS)) were conducted by transmission electron microscopy (Thermo Fisher, Talos F200X G2). EDS mapping data were processed with Velox software (Thermo Fisher). The crystal structures were carried out using a powder x-ray diffractometer (XRD) (Rigaku, D/MAX-RC 12 kW) using Cu $K\alpha$ ($\lambda = 1.54 \text{ \AA}$) radiation. The chemical elements and bonding states were analyzed using x-ray photoelectron spectroscopy (XPS) (Sigma Probe, Thermo VG Scientific) with Al $K\alpha$ radiation (1486.7 eV). Raman spectra was collected by a Horiba Jobin Yvon ARAMIS with a 514 nm excitation source. For the atomic-resolution imaging of a single atom on MoS_2 , high-angle annular dark field scanning transmission electron microscopy (HAADF-STEM) was conducted with double-Cs-corrected transmission electron microscope (Spectra Ultra, Thermo Fisher). The microscope operated at an accelerating voltage of 200 kV and a beam convergence angle of 21.4 mrad. HAADF-STEM images were obtained with a screen current of $\sim 20 \text{ pA}$, a resolution of 4096×4096 pixels, a pixel size of 6.249 pm, and a dwell time of 500 ns. For HAADF-STEM imaging, the detector's inner and outer collection angles were set to 52 and 200 mrad, respectively.

4.6 | Gas Sensing Measurements

The sensing materials, MoS_2 samples, were initially coated on the Al_2O_3 substrate using the drop-coating method. Each sensing material was dispersed in ethanol (2 mg in 100 μL) and drop-coated onto the substrate using a micropipette (2.5 μL) several

times at room temperature. Gas sensing tests were conducted in a homemade gas sensing system featuring a 16-channel multiplexer (34902A, Agilent), a data acquisition system (34972A, Agilent) that measured resistance across 16 channels at 4-s intervals, and a mass flow control system for regulating gas concentration (Figure S61). Sensing tests were conducted in a humid environment (0%, 50%, and 75% RH), controlled by a mixture of water bubbler and balancing gas with mass flow controllers. Sensitivity, selectivity, and durability tests involved cyclic on/off exposures (30 and 60 min) to various gases and air, with response calculated as the change in resistance ($\Delta R/R_{\text{air}}$ [%]), indicating the ratio of baseline resistance in the air (R_{air}) to resistance change after gas exposure. The baseline resistances were stabilized over 3 h at room temperature. The sensor substrates consist of three main components: an alumina board, gold electrodes, and a platinum heater. The alumina substrate measures 2.5 mm $\times$ 2.5 mm with a thickness of 0.2 mm and printed interdigitated Au electrodes (width: 25 μm , spacing: 70 μm). The ex situ Raman spectrometer was performed with the sensing materials after exposure for 1 h to 5 ppm of NO_2 .

4.7 | Computational Details

The spin-polarized DFT calculations were conducted using Vienna Ab-initio Simulation Package (VASP) with the projector-augmented wave (PAW) method [49–51]. Wave functions were represented using a plane-wave basis set with an energy cutoff of 520 eV. Perdew–Burke–Ernzerhof (PBE) functional was used to calculate the exchange-correlation energy [52]. DFT-D3 method with zero-damping function was employed to describe van der Waals interactions [53]. The convergence criterions for electronic and ionic step were set to 10^{-6} eV, and $-0.02 \text{ eV/\AA}$, respectively. The number of valence electrons was set to 5, 6, 6, 6, 9, 10, and 11 for N, O, Mo, S, Ir, Pt, and Au, respectively.

The FTS- MoS_2 slab model for adsorption energies and Bader charge differences was constructed with an interface between MoS_2 and Mo_2S_3 and a sulfur vacancy (V_s). $\text{MoS}_2/\text{Mo}_2\text{S}_3$ interface was generated by placing the MoS_2 (001) slab on the Mo_2S_3 (-101) slab with interface builder in QuantumATK [54] (Figure S19). To accommodate lattice mismatch, a mean absolute strain of 2.44% was introduced with directional strains $\epsilon_{11} = -2.41\%$, $\epsilon_{22} = 0.89\%$ and $\epsilon_{12} = -4.02\%$. A vacuum region of 15 $\text{\AA}$ was introduced along the z -axis. $1 \times 5 \times 1$ Γ -centered k -point grid was used to sample the Brillouin zone for these slab structures.

Adsorption energies of adsorbate (single atoms and NO_2) (E_{ads}) was calculated using the Equation 3:

$$E_{\text{ads}} = E_{\text{slab}^*} - (E_{\text{slab}} + E_{\text{adsorbate}}) \quad (3)$$

where E_{slab^*} , E_{slab} , and $E_{\text{adsorbate}}$ are total energies of slab model with adsorbate, bare slab and the isolated adsorbate, respectively. The energies of the isolated adsorbate were calculated in a $10 \times 10 \times 10 \text{ \AA}^3$ cubic cell with a $1 \times 1 \times 1$ Γ -centered k -point grid. VESTA software was utilized to illustrate crystal structures and charge density difference ($\Delta\rho$) maps for NO_2 adsorption were computed using the Equation 4:

$$\Delta\rho = \rho_{\text{slab}^*} - (\rho_{\text{slab}} + \rho_{\text{NO}_2}) \quad (4)$$

where ρ_{slab^*} , ρ_{slab} , and ρ_{NO_2} are charge density of NO_2 adsorbed slab, bare MoS_2 or $\text{M}_{\text{SA}}\text{-MoS}_2$ slab and NO_2 molecule. Isosurface level of charge density difference map was set to $0.01 \text{ e}/\text{\AA}^3$.

Author Contributions

I.-D.K. and D.-H.K. supervised the overall project. E.S. conceived the work, designed the whole project, including conceptualization and experiments, and wrote the manuscript. J.C. and J.M.Y. carried out S/TEM measurements. W.C., S.P., and C.-W.L. carried out DFT calculations. S.W. helped to carry out gas sensing measurements. M.S., M.S.K., C.P., and J.W.B. assisted the measurements and material synthesis. S.-Y.C. assisted the IPL process. All authors discussed the results and assisted during manuscript preparation.

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Conflicts of Interest

The authors declare no conflicts of interest.

Data Availability Statement

The data that support the findings of this study are available on request from the corresponding author. The data are not publicly available due to privacy or ethical restrictions.

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Supporting Information

Additional supporting information can be found online in the Supporting Information section.

Supporting File: adma75106-sup-0001-SuppMat.docx.