# Sensing Behavior of Atomically Thin-Layered MoS<sub>2</sub> Transistors

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**ABSTRACT** Most of recent research on layered chalcogenides is understandably focused on single atomic layers. However, it is unclear if single-layer units are the most ideal structures for enhanced gas—solid interactions. To probe this issue further, we have prepared large-area MoS<sub>2</sub> sheets ranging from single to multiple layers on 300 nm SiO<sub>2</sub>/Si substrates using the micromechanical exfoliation method. The thickness and layering of the sheets were identified by optical microscope, invoking recently reported specific optical color contrast, and further confirmed by AFM and Raman spectroscopy. The MoS<sub>2</sub> transistors with different thicknesses were assessed for gas-sensing performances with exposure to NO<sub>2</sub>, NH<sub>3</sub>, and humidity in different conditions such as gate bias and light irradiation. The results show that, compared to the single-layer counterpart, transistors of



few MoS<sub>2</sub> layers exhibit excellent sensitivity, recovery, and ability to be manipulated by gate bias and green light. Further, our *ab initio* DFT calculations on single-layer and bilayer MoS<sub>2</sub> show that the charge transfer is the reason for the decrease in resistance in the presence of applied field.

**KEYWORDS:** MoS<sub>2</sub> · gas sensor · gate bias · light irradiation · density functional theory

raphene has attracted much attention in recent years, with many potential proof-of-concept device demonstrations, such as field effect transistors,<sup>1–5</sup> ultracapacitors,<sup>6–8</sup> solar cells,<sup>9–12</sup> and especially gas sensors.<sup>13-15</sup> For gas sensors, the prior experience of traditional semiconducting metal oxides indicates improved sensing performance with an increase in the surface-to-volume ratio.<sup>16</sup> This is why graphene is believed to outperform traditional sensors since one atomically thin 2-D form has exceedingly high surface-to-volume ratio. Indeed, previous reports on graphene show excellent sensing performance.<sup>14</sup> However, in addition to surface-to-volume ratio, other important factors for a good gas sensor include the semiconducting properties<sup>17,18</sup> and availability of reactive sites for redox reactions. The semiconducting behavior is particularly attractive for modulating the transport characteristics with exposure to light or gate bias to enhance sensing performance.<sup>19,20</sup> Graphene fails in this regard because it has zero band gap. Recently, the lack of a band gap of graphene has inspired researchers to search for other graphene-like 2-D layered semiconducting materials with appropriate band gap. To date, there are several other atomically thin-layered 2-D materials that have been identified and investigated, such as molybdenum disulfide ( $MOS_2$ ),<sup>21–33</sup> gallium sulfide (GaS),<sup>26,34</sup> gallium selenide (GaSe),<sup>26,34</sup> hexagonal boron nitride (h-BN),<sup>35–40</sup> niobium diselenide ( $NDSe_2$ ),<sup>32,40–42</sup> tungsten diselenide ( $WSe_2$ ),<sup>43,44</sup> among others.

Among these emerging 2-D layered materials, MoS<sub>2</sub> is an n-type semiconductor and has been widely studied. More importantly, unlike graphene, the bulk MoS<sub>2</sub> possesses an indirect band gap (1.2 eV). Remarkably, for two-dimensional single-layer sheets, the band gap becomes direct and wider (1.8 eV). This transition from indirect band gap to direct band gap has also been confirmed by a strong photoluminescence of single-layer MoS<sub>2</sub>.<sup>45,46</sup> However, in spite of excellent semiconducting properties and the possibility of large surface-to-volume ratio of atomically thin-layered MoS<sub>2</sub>, there are few recent reports on its potential as a

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Figure 1. Single-layer MoS<sub>2</sub> sheet deposited on 300 nm SiO<sub>2</sub>/Si substrate using the micromechanical cleavage method. (a) Optical image, (b) Raman image, (c) AFM image, and (d) corresponding AFM height profile.

viable and effective chemical sensor.<sup>47–49</sup> Further, most of these reports do not emphasize the subtleties of gas-sensing performance and its relation to layered structures but instead demonstrate gas sensing as a proof-of-concept.

Herein, for the first time, we have prepared and wellcharacterized single- and multiple layers of crystalline MoS<sub>2</sub> sheets to fabricate a transistor-based gas-sensing platform. The sensing performance of these transistors for NO<sub>2</sub>, NH<sub>3</sub>, and humidity with and without gate bias and light irradiation has been assessed to unravel the role of layered structures in gas—solid interactions.

# **RESULTS AND DISCUSSION**

After the MoS<sub>2</sub> sheets were deposited by the mechanical exfoliation method, we used a complementary approach to confirm their thickness. Figure 1a shows an optical image of a single-layer MoS<sub>2</sub> sheet deposited on a 300 nm SiO<sub>2</sub>/Si substrate using the micromechanical cleavage method. Figure 1b shows Raman image of single-layer MoS<sub>2</sub> sheet. Figure 1c shows the AFM image, and Figure 1d shows the corresponding quantitative AFM height profile. The measured height of a single-layer MoS<sub>2</sub> sheet is ~0.9 nm (Figure 1d), which is consistent with previous reports.<sup>26,47</sup> For 2–5 layers of MoS<sub>2</sub>, the measured thicknesses were approximately 1.4, 2.2, 2.5, and 3.3 nm, respectively (as shown in Supporting Information Figures S12–S15).

Figure 2 shows TEM images of few-layer MoS<sub>2</sub> sheets transferred onto a QuantaFoil TEM grid: Figure 2a shows the low-magnification TEM image; the inset

shows the high-resolution TEM image with Mo–S spacing of 0.3-0.32 nm. Figure 2b is the high-resolution TEM image showing Moiré fringes due to the overlap of multiple MoS<sub>2</sub> layers; the distance between fringes is between 0.6 and 0.7 nm, which is comparable to a single layer of the S–Mo–S building block. Figure 2c shows the low-magnification TEM image of few-layer MoS<sub>2</sub> sheets. We have also noticed amorphization of MoS<sub>2</sub> layers due to high-energy electron radiation. The region is marked in a circle, as shown in Figure 2c. Figure 2d shows the selected area electron diffraction (SAED) pattern of MoS<sub>2</sub> layers, confirming the hexagonal structure and that the orientation of the MoS<sub>2</sub> sheet is along the  $\langle 100 \rangle$  zone.

Raman spectroscopy gives precise and exact information about properties of the material such as fine structure, atomic bonds, thermal expansion, specific heat, thermal conductivity, crystallographic orientation, stress, chemical composition, phase transition, etc. Further, the localized or extended (mapping) can also be studied using micro-Raman. It is also important to know the vibrational properties of single-layer to multilayer MoS<sub>2</sub> sheets to understand the electronphonon interaction and the transport properties of MoS<sub>2</sub> sheets, which have a large impact especially on field effect transistor device performance. Raman spectroscopy is quite sensitive to the thickness of 2-D layered material systems, and our results are consistent with the reported values.<sup>26,50</sup> Figure 3a shows the comparative Raman spectrum of bulk and single-layer MoS<sub>2</sub>. Figure 3b shows the Raman spectrum of the

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Figure 2. TEM images of few-layer MoS<sub>2</sub> sheets transferred onto a QuantaFoil grid. (a) Low-magnification TEM image; inset shows the high-resolution TEM image. (b) High-resolution TEM image showing Moiré fringes due to the overlap of multiple MoS<sub>2</sub> layers. (c) Low-magnification TEM showing electron-beam damage. (d) SAED pattern.

bilayer MoS<sub>2</sub> sheet before and after exposure to 1000 ppm NH<sub>3</sub> gas using 514 nm Ar laser irradiation. The position of  $E_{2q}^{1}$  and  $A_{1q}$  mode stiffens and broadens due to exposure of 1000 ppm NH<sub>3</sub> gas with bilayer MoS<sub>2</sub> sheets. Our results of the Raman spectrum on exposure to NH<sub>3</sub> gas is consistent with a recent report of charge transfer interaction with an electron donor molecule.<sup>51,52</sup> Figure 3c shows the effect of laser power variation on the peak position of the Raman  $E_{2q}^{1}$  mode for single-layer to six layers of MoS<sub>2</sub>. Figure 3d shows the effect of laser power variation on FWHM of the  $E_{2q}^{1}$ mode for single-layer to six layers of MoS<sub>2</sub>. Figure 3e shows the effect of laser power variation on peak position of the A1g mode for single-layer to six layers of MoS<sub>2</sub>, and Figure 3f shows the effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer MoS<sub>2</sub> to six layers of MoS<sub>2</sub> sheet samples. Single-layer MoS<sub>2</sub> to multilayer MoS<sub>2</sub> samples show a red shift for both  $E_{2q}^{1}$  and  $A_{1q}$  modes. Further, we observed the increased line width as a function of increasing laser power. The change in peak position as well as peak broadening as a function of increasing laser power indicates the increasing local temperature on the MoS<sub>2</sub> sheet sample surface. Thus, the Raman spectrum of MoS<sub>2</sub> as a function of number of layers shows laser power dependence. The  $E_{2g}^1$  and  $A_{1g}$  are sensitive to laser power variation, and change in peak position and line width of the spectra was observed.<sup>26,53</sup>

The gas-sensing device platform was fabricated using electron-beam lithography (E-BL) for patterning Ti/Au electrodes on top of isolated MoS<sub>2</sub> sheets, and back gate geometry was used. Figure 4a shows the schematic of the MoS<sub>2</sub> transistor-based NO<sub>2</sub> gassensing device, and Figure 4b shows the optical photograph of the sensing device mounted on the chip. The electrodes on the sample were connected to the gold pad on the chip carrier by gold wire. This chip carrier was then mounted on the stage inside the gas testing chamber for measurement. Figure 4c shows the optical image, and Figure 4d shows the SEM image of an E-beam lithography patterned two-layer MoS<sub>2</sub> transistor device. In a typical MoS<sub>2</sub> transistor device, the length of the sheet is about 30  $\mu$ m, and the spacing of the electrodes is 5  $\mu$ m. Extra efforts were spent during the mechanical exfoliation process to deposit large enough flakes for electrode design, which not only facilitates patterning electrodes but also helps diminish edge portion since edge sites may act differently in the gas-sensing experiment. Figure 4e shows the experimental setup used for NH<sub>3</sub> and NO<sub>2</sub> gas-sensing performance. The as-fabricated singleto five-layer MoS<sub>2</sub> FETs were tested with different gases at room temperature. We found that the current of the single-layer MoS<sub>2</sub> device was not stable over time as reported earlier,47 while multilayer MoS<sub>2</sub> devices exhibited much more stable response. For clarity and brevity, we will only discuss two-layer and five-layer samples in this article because they are the thinnest and the thickest. We note that both three- and four-layer samples seem to show behavior closer to that of the fivelayer sample.

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Figure 3. (a) Raman spectrum of bulk and single-layer  $MoS_2$  sheet. (b) Raman spectrum of bilayer  $MoS_2$  sheet before and after exposure to 1000 ppm  $NH_3$ . (c) Effect of laser power variation on peak position of the Raman  $E^{1}_{2g}$  mode for single-layer to six layers of  $MoS_2$ . (d) Effect of laser power variation on the FWHM of the  $E^{1}_{2g}$  mode for single-layer to six layers of  $MoS_2$ . (e) Effect of laser power variation on the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$ . (f) Effect of laser power variation on the FWHM of the A<sub>1g</sub> mode for single-layer to six layers of  $MoS_2$  sheets.

Figure 5 shows the gas-sensing behavior of twolayer and five-layer MoS<sub>2</sub> FETs when exposed to NH<sub>3</sub> (Figure 5a) and NO<sub>2</sub> (Figure 5b). The analyte gas was injected into the chamber for 3 min, and then samples were allowed to recover for 10 min in pure nitrogen environment. There were four cycles in this experiment, and the analyte gas concentration in each cycle was chosen to be 100, 200, 500, and 1000 ppm. Sensitivity is defined as  $S = (R_q - R_{N_2})/R_{N_2}$ , where  $R_q$  is the resistance of the device in target gas, and  $R_{N_2}$  is the resistance of the device in dry N<sub>2</sub>. It is well-known that NO<sub>2</sub> acts as an electron acceptor (p-type doping) while NH<sub>3</sub> acts as an electron donor (n-type doping). Thus, for n-type semiconductors, such as MoS<sub>2</sub>, NO<sub>2</sub> increases its resistance while NH<sub>3</sub> decreases its resistance based on a charge transfer mechanism. Our results show clear response of MoS<sub>2</sub> layers to both NO<sub>2</sub> and NH<sub>3</sub>, and the trends in resistance change confirm that the charge transfer mechanism

dominates the sensing process. The recovery is often not fully complete in each cycle, which is also observed in graphene-based sensors<sup>15</sup> and in previous MoS<sub>2</sub> sensor reports.<sup>47</sup>

Interestingly, we notice that the five-layer  $MoS_2$  sample has better sensitivity compared to that of the two-layer  $MoS_2$  sample, which seems counterintuitive. We should point out that samples were tested multiple times with more than 1 day interval between the tests to allow full recovery of test samples, but the results remained the same. This phenomenon may be due to the different electronic structures when thicknesses (layering) are dissimilar. The energy levels of conduction and valence band of  $MoS_2$  changes with size, as shown in an earlier report,<sup>54</sup> and thus they may likely alter the redox potentials, which in turn affect the sensing ability of  $MoS_2$ . Clearly, this is a complex issue and still under study to clarify the specific sensing mechanism.

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Figure 4. (a) Schematic of the  $MoS_2$  transistor-based  $NO_2$  gas-sensing device. (b) Optical photograph of the  $MoS_2$  sensing device mounted on the chip. (c) Optical image of two-layer  $MoS_2$  sheet. (d) SEM image of two-layer  $MoS_2$  transistor device, and (e) experimental setup used for  $NH_3$  and  $NO_2$  gas-sensing performance.



Figure 5. Comparative two- and five-layer MoS<sub>2</sub> cyclic sensing performances with (a) NH<sub>3</sub> and (b) NO<sub>2</sub> (for 100, 200, 500, 1000 ppm).

In addition to the cyclic experiments shown in Figure 5, we also have carried out measurements of  $I_{ds}$  (drain-source current)– $V_{gs}$  (gate-source voltage) curves with fixed  $V_{ds}$  (drain-source voltage) before and after exposure to NH<sub>3</sub> and NO<sub>2</sub>, as shown in Supporting Information Figure S3. This shows that the semiconducting property of MoS<sub>2</sub> sheets can be tuned by exposure to gases. Figure 6 shows the comparative sensing behavior with and without applying back gate voltage (+15 V) for two-layer NH<sub>3</sub>

(Figure 6a) and NO<sub>2</sub> (Figure 6b) and five-layer MoS<sub>2</sub> NH<sub>3</sub> (Figure 6c) and NO<sub>2</sub> (Figure 6d). As mentioned in the introduction, resistance in MoS<sub>2</sub> can be tuned by gate biasing, which makes MoS<sub>2</sub> more competitive for gas sensing compared to graphene-based sensors. Interestingly, applying back gate voltage affects gassensing responses differently for reducing gas (NH<sub>3</sub>) and oxidizing gas (NO<sub>2</sub>). Applying positive back gate voltage enhances sensitivity for NO<sub>2</sub>, while decreasing sensitivity for NH<sub>3</sub>. The phenomena can be understood

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Figure 6. Comparative sensing behavior with and without applying back gate voltage (+15 V) for two-layer (a)  $NH_3$ , (b)  $NO_2$  and five-layer  $MOS_2$  (c)  $NH_3$  and (d)  $NO_2$  (for 100, 200, 500, 1000 ppm).



Figure 7. Sensitivity as a function of concentration (in ppm) for two-layer and five-layer MoS<sub>2</sub> sheets for (a) NH<sub>3</sub> and (b) NO<sub>2</sub>.

as follows: when positive gate voltage is applied, electrons accumulate at the interface of MoS<sub>2</sub> and SiO<sub>2</sub>, so for NO<sub>2</sub>, which is an electron acceptor, it has more electrons to accept. However, for NH<sub>3</sub>, which is an electron donor, the electric field built up at the interface repels NH<sub>3</sub> to donate an electron, which in turn decreases the gas sensitivity. Furthermore, by comparing the change in sensitivity with and without gate bias, the five-layer MoS<sub>2</sub> sample showed stronger response to gate bias, which suggests that carrier concentration in five-layer MoS<sub>2</sub> is more susceptible to the influence of gate bias. This can be further confirmed by the higher ON/OFF ratio of the five-layer sample shown in Supporting Information Figure S3. The ON/OFF ratio of five-layer MoS<sub>2</sub> is  $5.5 \times 10^5$ , compared to 7.5 × 10<sup>2</sup> of two-layer MoS<sub>2</sub>. This gatedependent selectivity can be effectively utilized to make MoS<sub>2</sub> sensors more selective to certain type of gases, as this is a much easier approach than incorporating dopants into material which is the common way to make sensors selective to certain test gas.<sup>55</sup> Figure 7 shows the summary of sensitivity as a function of concentration (in parts per million) for two-layer and five-layer MoS<sub>2</sub> sensor devices in different operational situations with exposure to (a) NH<sub>3</sub> and (b) NO<sub>2</sub>. The highest sensitivity toward NO<sub>2</sub> was achieved on the five-layer MoS<sub>2</sub> sample operated with positive gate voltage. The highest sensitivity toward NH<sub>3</sub> was achieved on the five-layer MoS<sub>2</sub> sample operated without gate voltage. Their sensitivities when exposed to

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Figure 8. (a) Photoluminescence spectra of five-layer  $MoS_2$ . (b) Change of normalized resistance under light irradiation. (c) Sensing behavior of five-layer  $MoS_2$  exposed to 100 ppm  $NO_2$  under green light illumination. (d) Sensing behavior of five-layer  $MoS_2$  exposed to 100 ppm  $NH_3$  under green light illumination.

1000 ppm  $NO_2$  and  $NH_3$  were determined to be 1372 and 86%, respectively.

Other than applying gate bias, light exposure is long believed to help increase sensitivity and shorten recovery time<sup>19</sup> attributed to varied mechanisms. We applied the same approach to the five-layer MoS<sub>2</sub> sample for a series of light-assisted sensing experiments. Figure 8a shows the photoluminescence (PL) spectrum of the five-layer MoS<sub>2</sub> sample with a peak located at 687 nm, confirming the existence of finite band gap in the material, and the result is similar to the reported value.<sup>45,46</sup> To further examine the influence of light irradiation, we then exposed the sample to green LED (532 nm) whose power density can be tuned by changing currents running through it and recorded the variation of resistance. From Figure 8b, it can be seen that resistance goes down when the sample is irradiated by light, and the magnitude depends on the illuminated power density (the three curves are deliberately shifted horizontally for clarity).

The influence of light can be divided into two stages: at first stage, resistance drops quickly with a small amount, and at second stage, resistance slowly reaches saturation with a larger amount. From a recently published report,<sup>56</sup> the photocurrent generated in monolayer  $MoS_2$  is attributed mostly to heat rather than exciton, which explains in our  $MoS_2$  sensor device that the second stage of photocurrent is more dominant than the first stage, especially that five-layer  $MoS_2$ is an indirect semiconductor while monolayer  $MoS_2$ is a direct semiconductor. Figure 8c,d shows the gas-sensing behavior of the five-layer MoS<sub>2</sub> sample reacting with NO<sub>2</sub> and NH<sub>3</sub>, respectively, under green light irradiation with different power densities. Although many similar light-assisted sensing experiments use UV light as the source, the damage of UV light needs to be taken into account, as shown in a report of carbon-based sensors.<sup>57</sup> Furthermore, since the device showed response to areen light, we decided to use it to minimize damage. Further, the SEM and optical microscope were used to observe the effect of green light irradiation, which showed no obvious and apparent damage caused after long-time green light exposure. In order to compare the response and recovery, we injected analyte gas for a longer time (20 min) and also let the device recover for a longer time (30 min). For NO<sub>2</sub>, the sensitivity increased when exposed to 4 mW/cm<sup>2</sup> irradiation but decreased when exposed to 15 mW/cm<sup>2</sup>. It is similar to the behavior of UV-assisted metal oxide sensors, for which the power density of light has an optimal value.<sup>58</sup> There are two possible reasons; one is that when the power density is too high, not all the excited electrons/holes get to react with gas molecules, so the portion of reacted carriers decreases; and the other possible reason is that the desorption rate increases more than the adsorption rate because of the light-induced activation under irradiation. For response and recovery, it is obvious that the higher the power density, the quicker the response/recovery. However, in the case of NH<sub>3</sub>, as shown in Figure 8d, the behavior is different. Although the sensitivity still has an optimal value, where

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Figure 9. Current–voltage (I-V) behaviors measured under different RH values, showing an obvious slope decreasing with the increasing of RH for (a) two-layer and (b) five-layer MoS<sub>2</sub> sample. RH as a function of resistance for (c) two-layer and (d) five-layer MoS<sub>2</sub> samples. Inset of (c) shows response and recovery for change in RH from dry air (4%) to 84% RH. Response time is quick, but recovery is longer for MoS<sub>2</sub> layered material samples.

4 mW/cm<sup>2</sup> gave the highest sensitivity, the response and recovery did not show a difference whether the sample was exposed to light or not. Higher power density was also tried as shown in Supporting Information Figure S4, which indicated the response/recovery behavior still does not change. To confirm whether NH<sub>3</sub> is the special case, we have also performed similar experiment on H<sub>2</sub>S, as shown in Supporting Information Figure S5, and there is no clear difference from NH<sub>3</sub>. It implies that the kinetic energy barrier for reducing gas molecules to leave the MoS<sub>2</sub> surface is likely higher than that of NO<sub>2</sub>. To our knowledge, this is the first report that light irradiation does not appreciably assist in response/recovery, and further studies are warranted to shed further light on this observation.

The humidity sensing performance of two-layer and five-layer MoS<sub>2</sub> sensor devices was also investigated to complete the sensing portfolio. A recent study concluded that a single-layer MoS<sub>2</sub> device does not respond to water vapor.<sup>49</sup> However, another previous work<sup>54</sup> showed that nano-MoS<sub>2</sub> reacted with water, and our own prior results also showed that the currentvoltage (I-V) behavior of MoS<sub>2</sub> changed in the presence of water vapor.<sup>30,31</sup> In this work, we measured the I-V behavior under different relative humidity (RH) values, and results showed obvious decrease in the slope of I-V curves with an increase in RH, as shown in Figure 9a for the two-layer MoS<sub>2</sub> sensor device and Figure 9b for the five-layer MoS<sub>2</sub> sensor device. Water vapor is considered to be an electron acceptor just like NO<sub>2</sub>, so the resistance of MoS<sub>2</sub> increases with the RH

value. Once again, this observation further confirms that charge transfer mechanism is dominant in this process. Figure 9c,d shows the resistances as functions of RH of two-layer  $MoS_2$  and five-layer  $MoS_2$  sensor devices, respectively. The inset of Figure 9c shows response and recovery for change in RH from dry air (4%) to 84% RH.

To understand the experimental results on the sensing behavior of thin-layered  $MoS_2$  transistors, we have carried out the *ab initio* calculations on single- and bilayer  $MoS_2$ . A 3  $\times$  3 supercell of  $MoS_2$  was constructed to study the adsorption of  $NO_2$  and  $NH_3$  on both single-layer (ABA stacking of intralayer S-Mo-Satoms) and bilayer (ABA, BAB stacking sequence of the adjacent layers)  $MoS_2$ .<sup>59</sup> We have computed energy of adsorption for both  $NH_3$  and  $NO_2$  gases on mono- and bilayer  $MoS_2$  using the following expression

$$E_{ads}[MoS_2 + X] = E_T[MoS_2 + X] - E_T[MoS_2] - E_T[X]$$

where  $E_{T}[MOS_{2} + X]$  is the total energy of the supercell and molecule,  $E_{T}[MOS_{2}]$  is the total energy of the MOS<sub>2</sub> slab, and  $E_{T}[X]$  is the energy of the adsorbed molecule. The adsorption energy for both NH<sub>3</sub> and NO<sub>2</sub> on singleand bilayer MOS<sub>2</sub> was estimated in the presence and absence of external electric field of 0.0001 au (which is equivalent to  $V_{g} = 15$  V). Our results show that there is a very slight increase in the adsorption with an increase in the number of MOS<sub>2</sub> layers (*i.e.*, from single-layer to bilayer), hence indicating very small increase in sensitivity for five layers as compared to two layers. The adsorption energy of NO<sub>2</sub> is greater than that of NH<sub>3</sub>

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TABLE 1. Comparison of Adsorption Energy for Different Supercell Configurations (kJ/mol)



Figure 10. (a) Charge transfer in NH<sub>3</sub> on bilayer MoS<sub>2</sub> with  $V_g = 0$  V. The red isosurface indicates negative charge, and blue isosurface denotes positive charge. (b) Variation of resistance for NH<sub>3</sub> on bilayer MoS<sub>2</sub> with  $V_g = 0$  and 15 V.

 TABLE 2. Comparison of Charge Transfer for Different

 Configurations (Electrons/Cell)<sup>a</sup>

supercell gas bilayer MoS <sub>2</sub>	$V_{g} = 0 V$		$V_{\rm g}=15~{ m V}$	
	NH3 -0.009	NO <sub>2</sub> +0.097	NH <sub>3</sub> 0.018	NO <sub>2</sub> +0.089

<sup>*a*</sup> Negative is for charge transferred to MoS<sub>2</sub>, and positive is for charge transferred from MoS<sub>2</sub>.

for both single-layer and bilayer  $MoS_2$ , as shown in Table 1. The higher the adsorption energy, the higher its binding with the substrate, and hence  $MoS_2$  is more sensitive to  $NO_2$  than  $NH_3$ .

We have also estimated the amount of charge transferred between MoS<sub>2</sub> and NH<sub>3</sub>, NO<sub>2</sub> gases in the presence and in the absence of external electric field. MoS<sub>2</sub> is an n-type semiconductor, implying an excess of electrons and hence conducting. NH<sub>3</sub> acts as an electron donor and donates electrons when absorbed on MoS<sub>2</sub>, as shown in Figure 10a. In the presence of a positive gate voltage, electrons in the NH<sub>3</sub> molecule are attracted to the MoS<sub>2</sub>, which leads to more electrons transferred to MoS<sub>2</sub> as compared to the number of electrons transferred in the absence of an electric field, as shown in Table 2. Due to the increase in charge carrier resistance, the MoS<sub>2</sub> sheet decreases and is in agreement with experimental observations, as shown in Figure 10b. With the increase in the concentration of NH<sub>3</sub>, the charge transferred to the MoS<sub>2</sub> sheet increases. Due to the increase in repulsive energy between the transferred electrons and incipient electrons in MoS<sub>2</sub>, further transfer of electrons slows with



Figure 11. (a) Charge transfer in NO<sub>2</sub> on bilayer MoS<sub>2</sub> with  $V_g = 0$  V. The red isosurface indicates negative charge, and blue isosurface denotes positive charge. (b) Variation of resistance for NO<sub>2</sub> on bilayer MoS<sub>2</sub> with  $V_g = 0$  and 15 V.

concentration of gas. Hence, the resistance of the  $MoS_2$  sheet decreases at first significantly and decreases slowly after a certain concentration, as shown in Figure 10b.

On the other hand, NO<sub>2</sub> acts as an electron acceptor and accepts electrons from the MoS<sub>2</sub> sheet, as shown in Figure 11a. This charge transfer reduces the number of electrons in the MoS<sub>2</sub> sheet, thus reducing the charge carriers in it. On application of a positive gate voltage, the electrons in the MoS<sub>2</sub> sheet are attracted toward the gate, thus reducing the charge transferred to NO<sub>2</sub>, as shown in Table 2. Hence, the resistance of the substrate decreases with the application of external field. With the increase in the concentration of  $NO_{2}$ , more charges are transferred from MoS<sub>2</sub> to NO<sub>2</sub>. The charges are transferred to different NO<sub>2</sub> molecules as the concentration increases, and unlike NH<sub>3</sub>, there is no repulsion between these charges. Hence, resistance of the MoS<sub>2</sub> sheet steeply increases with the concentration of NO<sub>2</sub>, as shown in Figure 11b. Our first-principles calculations are hence in agreement with experimental observations.

Our calculations explain the changes in overall resistance of the MoS<sub>2</sub> channel with gas adsorption and gate voltage. However, to understand the trends in the relative resistance with gate voltage, we find that the resistance of the contact of MoS<sub>2</sub> with Ti/Au electrodes is also relevant in addition to the charge transfer mechanism. The variation of electrostatic potential of the bulk  $MoS_2$  along the z direction (along the direction perpendicular to the MoS<sub>2</sub> planes) with and without electron doping shows a high barrier of  $\approx$ 11.56 eV between the MoS<sub>2</sub> layers, as shown in Figure 12. This implies that the electrons have to cross over a large barrier for movement toward the gate and that electrons mostly from the topmost single-layer and its contact with source and drain electrodes determine the electrical properties of the transistor. In the presence of positive gate voltage, the electrons of the Ti/Au electrode accumulate at the MoS<sub>2</sub> and Ti/Au contacts at the source and drain, reducing the

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Figure 12. Variation of macroscopic average of the Kohn–Sham potential for bulk  $MoS_2$  with and without electron doping along the *z* direction.

barrier between the Ti/Au electrodes and the  $MoS_2$ sheet and easing the transfer of electrons across the contacts. Hence, even in presence of the ambient  $N_2$ gas (zero concentration of  $NH_3$  or  $NO_2$ ), a large reduction in the resistance is observed with the application of gate voltage, which is responsible for the observed trends in the relative changes in the channel resistance.

## CONCLUSION

We report a comprehensive suite of sensing behavior of  $MoS_2$  atomically thin-layered structures in a transistor geometry. Isolated MoS<sub>2</sub> sheets were prepared on 300 nm SiO<sub>2</sub>/Si substrates using micromechanical exfoliation and characterized for their thickness and layering with an optical microscopy, AFM, and Raman spectroscopy. Surprisingly, compared to the bilayer MoS<sub>2</sub> device, multilayer MoS<sub>2</sub> FETs showed more sensitive response to NH<sub>3</sub>, NO<sub>2</sub>, and also water vapor at room temperature and atmospheric pressure. We also successfully manipulated sensing behavior of MoS<sub>2</sub> by applying either gate bias or light exposure, which provide opportunities to improve sensing behavior and make MoS<sub>2</sub> to selectively respond to certain kinds of gases. Although the exact sensing mechanism of MoS<sub>2</sub> behavior is unclear and clearly has subtleties and nuances, it shows excellent sensing properties, worthy of further investigation. Our first-principles calculation shows that the charge transfer between the molecule and MoS<sub>2</sub> sheet channel is responsible for the decrease in resistance with applied field and suggests that the effect of contact resistance is also relevant to the observed changes in relative resistance. The contact resistance between the electrodes and MoS<sub>2</sub> layers is also an important underlying factor for the decrease in sheet resistance in the absence of absorbed gas (in the presence of  $N_2$ ).

# ARTICLE

## **EXPERIMENTAL METHODS**

**Marker on Si0<sub>2</sub>/Si Substrate.** For making markers, thermal evaporation was used to deposit Cr/Au (1 nm/50 nm) on a 300 nm SiO<sub>2</sub>/Si wafer with 200  $\mu$ m × 200  $\mu$ m squares with a separation of 1.5 mm (as shown in Supporting Information Figures S34 and S35) using shadow mask purchased from Photo Etch Technology, Lowell, MA (USA). Then substrates were cut into 10 mm × 10 nm square pieces using a diamond cutter pen purchased from SPI (USA) followed by standard substrate cleaning.

**SiO<sub>2</sub>/Si Substrate Cleaning.** First, SiO<sub>2</sub>/Si substrates were thoroughly cleaned by boiling substrates in acetone for 5 min and then boiling in IPA for 5 min followed by drying using dry N<sub>2</sub> gas. Finally, substrates were cleaned using O<sub>2</sub> plasma cleaning for 1 min.

Preparation of MoS<sub>2</sub> Sheets Using the Micromechanical Exfoliation Method. Bulk MoS<sub>2</sub> crystal (purchased from SPI Supplies Inc., USA, crystal size small, #429MM-AB 15 mm  $\times$  9 mm). MoS<sub>2</sub> crystals were collected by a professional geologist at Otter Lake deposit in Ontario, Canada. Single-layer to few-layer MoS<sub>2</sub> sheets were deposited at room temperature in ambient conditions by mechanically exfoliating the bulk MoS<sub>2</sub> crystal onto precleaned 300 nm SiO<sub>2</sub>/Si. First, the Scotch-tape was put on the MoS<sub>2</sub> bulk crystals, and then we repeatedly peeled off MoS<sub>2</sub> to find thin flakes. The next step was to transfer these thin MoS<sub>2</sub> flakes onto the SiO<sub>2</sub>/Si substrate by pressing the tape onto the substrates. Special care has been taken to exfoliate large-area MoS<sub>2</sub> sheets. First, the Scotch-tape was lightly pressed onto the MoS<sub>2</sub> bulk crystal and then removed carefully with a small angle as possible. Then peeling of the MoS<sub>2</sub> with Scotch-tape was done. We repeatedly peeled the MoS<sub>2</sub> on Scotch tape as long as we observed the light green color of MoS<sub>2</sub> flakes on the Scotch tape. The next step we followed was to transfer this light green color of MoS<sub>2</sub> flakes on Scotch tape onto the SiO<sub>2</sub>/Si substrate. We used plastic tweezers to put pressure on the Scotch tape horizontally from one side to the other side of the SiO<sub>2</sub> substrate for 30–60 s. This technique gives individual tens of micrometers

to 50  $\times$  50  $\mu m$  single-layer and few-layer MoS<sub>2</sub> sheets. Contaminants from the Scotch tape were removed by dipping the substrates in acetone for 5 min.

**Optical Microscopy.** The optical microscopic images were acquired with an optical microscope (Nikon Eclipse ME600) imager M1m with white light illumination (100 W halogen lamp, HAL100) using bright-field imaging modes and 100× objectives as described earlier.<sup>26</sup>

Atomic Force Microscopy. AFM tapping mode images were recorded using an ICON system (Bruker, Santa Barbara, CA, USA).

**Raman Spectroscopy and Raman Imaging.** Raman spectra and Raman imaging were done in back scattering geometry using the Nanophoton Raman-11 system with a 532 nm line of an Ar ion laser with laser power  $\sim 1$  mW as described earlier.<sup>26</sup>

**Raman Spectroscopy in NH<sub>3</sub> Gas.** Raman spectra of  $MoS_2$  in the presence of  $NH_3$  has been observed to prove n-type doping in a glass chamber (as shown in Supporting Information Figure S33) that consists of a gas inlet and outlet. To avoid the effect of ambient oxygen, which can cause the Raman shift in vibration modes by attacking the Mo-S-Mo bands, initially Ar gas was passed in test chamber to stabilize the  $MoS_2$  sample for 30 min. Then 1000 ppm NH<sub>3</sub> gas was passed in the test chamber. Raman spectra were recorded before and after the test gas was injected in the chamber.

**MoS<sub>2</sub> FET-Based Gas Sensor Device Fabrication.** MoS<sub>2</sub> FET devices were fabricated using electron-beam lithography as described earlier.<sup>30,31,34</sup> Electrodes were patterned by electron-beam lithography (EBL) followed by thermal evaporation to deposit titanium/gold (3 nm/70 nm). After metal deposition, samples were annealed in Ar atmosphere at 200 °C for 2 h to improve contact. For the back gate electrode, we simply scratched the surface and put indium metal as the contact. Samples were then wire bonded to the chip by gold wire to connect to the gassensing setup.

 $NO_2$  and  $NH_3$  Gas Sensor Test. Gas-sensing experiments were carried out in a home-built gas-sensing experimental setup,

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AGNANC www.acsnano.org as shown in Figure 4e. Gases we used for the test were NO<sub>2</sub> (2% NO<sub>2</sub> with 98% N<sub>2</sub>) and NH<sub>3</sub> (1% NH<sub>3</sub> balances 99% N<sub>2</sub>) cylinders purchased from Airgas Company (Port Allen, LA, USA). They were then mixed with pure nitrogen to achieve desirable concentrations before being injected into chamber. Mass flow controllers were used to control the concentration, and the total flow rate was 1000 sccm for the whole time. The testing chamber was a 250 mL flask, and gas should fill the chamber within 30 s based on the flow rate we used. The gas inlet was designed to point directly onto the sample with a distance of  $\sim$ 5 mm. Such a setup is for quick response to the change in environment. All the gas-sensing experiments were carried out at room temperature and at room atmospheric pressure.

**Green Light Illumination.** Green light illumination was carried out by using a LED head purchased from Sunlite Company (Lawrence, KS, USA). It was connected to a Keithley 2420 semiconductor to provide power, and the power density can be tuned by changing the current running through it. Electrical characterization and resistance change during gas injection were recorded with a Keithley 4200 semiconductor characterization system.

**Humidity Sensor Test.** For the humidity test, humidity was obtained using saturated salt solutions of various salts as described earlier by us for the graphene sensor.<sup>13</sup>

**Electrical Characterization.** Electrical characterization of field effect transistor-based sensing devices was performed by using a Keithley 4200 semiconductor characterization system to record the  $l_{ds}$  (drain-source current) change in real time (data acquisition time  $\sim$ 1 s) using the two-probe in a glass chamber. During the gas-sensing test, the MoS<sub>2</sub> sensor devices were always stabilized in nitrogen before testing, which usually took 30 min. All gate-bias-assisted and light-assisted tests were done on the same MoS<sub>2</sub> device to eliminate inaccuracy from device was stored in a vacuum box to speed recovery, and we made sure the MoS<sub>2</sub> sensor devices were identical.

Computational Details. Our first-principles calculations are based on density functional theory (DFT) with a generalized gradient approximated (GGA) exchange correlation energy functional of Perdew-Burke-Ernzerhof (PBE)<sup>60</sup> as implemented in the Quantum Espresso code.<sup>61</sup> Ultrasoft pseudopotentials<sup>62</sup> were used to capture the interaction between ionic cores and valence electrons. We use an energy cutoff of 30 Ry for the plane-wave basis used for the wave functions and 180 Ry to represent the charge density. We have included semiempirical dispersive van der Waals (vdW) interaction using Grimme parametrization.  $^{63}$  A 3  $\times$  3  $\times$  1 supercell was used to simulate the adsorption of NH<sub>3</sub> and NO<sub>2</sub> molecules with 1 molecule per supercell. Due to periodic boundary conditions, this geometry corresponds to 6.27  $\times$   $10^{13}$  molecules/cm². Integrations over the Brillouin zone were carried out using a regular mesh of 3 imes $3 \times 1$  k-points (for the structural relaxation) of the supercell and at the zone center for a single molecule. We have increased a denser mesh of  $6 \times 6 \times 1$  k-points for the charge transfer calculation. The amplitude of applied electric field was 0.0001 au, which corresponds to the back gate voltage ( $V_{\alpha}$ ) of +15 V applied in the experiments. The system was structurally relaxed until the maximum Hellmann-Feynmann force acting on each atom was less than 0.001 Ry/Å (for both, with and without electric field).

*Conflict of Interest:* The authors declare no competing financial interest.

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